

The chemomechanical constitutive equation for the poroelastic medium has the following form

$$\underline{\underline{\sigma}}' = \lambda \operatorname{tr} \underline{\underline{\varepsilon}}' \underline{\underline{I}} + 2\mu \underline{\underline{\varepsilon}}' \quad (10)$$

where the effective stress  $\underline{\underline{\sigma}}'$  tensor is defined ~~as~~ by the relation

$$\underline{\underline{\sigma}}' = \underline{\underline{\sigma}} + p_w \underline{\underline{I}} \quad (11)$$

and the chemoelastic strain  $\underline{\underline{\varepsilon}}'$  is defined by the relation

$$\underline{\underline{\varepsilon}}' = \underline{\underline{\varepsilon}} + \alpha_c C \underline{\underline{I}} \quad (12)$$

where  $\alpha_c$  is the coefficient of chemical contraction and  $\underline{\underline{\varepsilon}}$  is the <sup>total</sup> strain of the medium,  $\lambda$  and  $\mu$  are the Lamé constants of the solid skeleton.

The balance of momentum equations (3) and (4), lead to the following equivalent equations

$$\operatorname{div} \underline{\underline{\sigma}} = 0 \quad (13)$$

$$n^w (\vec{\nabla}_w - \vec{\nabla}_s) = -k_H \vec{\nabla} p_w \quad (14)$$

where  $k_H$  is the hydraulic permeability of the fluid diffusing within the pores of the solid.

Boundary and initial conditions for the uniaxial deformation of the strip, along the  $x = x_1$  direction, where we have the smallest dimension  $L_1$ , in the orthogonal strip of dimensions  $L_1 \times L_2 \times L_3$ .

Remember that  $L_1 \ll L_2 \ll L_3$ .

1) Initial condition for the displacement  $u_s$ , for  $x \in [-l, l]$  where  $L_2 = 2L$

$$(IC)_u: u_s(x, t=0) = 0 \quad (14)$$

2) Boundary Condition for  $u_s$

$$(BC)_u: u_s(x=0, t) = 0 \quad (15)$$

3) Boundary condition for the stress  $\sigma'_{xx}$

$$\sigma'_{xx}(x = \pm l, t) = \overbrace{\left( H \frac{\partial u_s}{\partial x} + 3K\alpha_c C \right)}^{\sigma'_{xx}}(x = \pm l, t) = 0 \quad (16)$$

4) Boundary condition for the pore pressure  $p_w$

$$(BC)_{p_w}; p_w(x = \pm l, t) = 0. \quad (15)$$

The diffusion problem of the chemical, uncouples from the deformation problem of the medium.

$$(FE)_c; \frac{\partial C}{\partial t} - D \frac{\partial^2 C}{\partial x^2} = 0 \quad (16)$$

$$(IC)_c; C(x, t=0) = 0 \quad (17)$$

$$(BC)_c; C(\pm l, t) = C_0 H(t) \quad (18)$$

with  $H(t)$  a step function

Returning to the deformation problem, we can combine the FE (mass and momentum balance equations) and end up with 2 equations in terms of 2 unknowns, the solid skeleton displacement  $u_s$  and the pore pressure  $p_w$ . These two equations are:

$$k_H \frac{\partial p_w}{\partial x} = \frac{\partial u_s}{\partial t} \quad (19)$$

$$\frac{\partial^2 u_s}{\partial x^2} - \frac{1}{HK_H} \frac{\partial u_s}{\partial t} + 3\alpha_c \frac{K}{H} \frac{\partial C}{\partial x} = 0 \quad (20)$$

and can be solved for  $u_s$  and  $p_w$ .

The tensile force  $F_z(t)$ , along the  $x_1$  longitudinal direction of the strip, is given via the relation



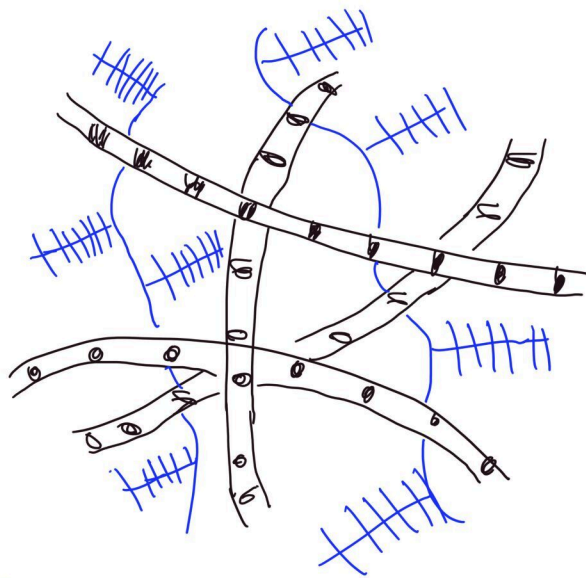
$$F(t) = L_2 \int_{-l}^l \sigma_{33}^s(x, t) dx \quad (21)$$

where

$$\sigma_{33}^s = \lambda \frac{\partial u_s}{\partial x} + 3K\alpha_c C - n^s p_w \quad (22)$$

elastic + chemical + fluid contributions to the stress

Illustration of the cartilage structure



Collagen fibrils

Proteoglycans

Proteoglycan repulsion ~~causes~~ force the collagen fibrils to stretch. Crucial role in the models for cartilage, play the electrochemical potentials:

$$g_k^{ec} = \underbrace{\hat{v}_k p_w}_{\text{mechanical part}} + \underbrace{RT L_k x_k}_{\text{chemical part}} + \underbrace{\sum_k F \phi}_{\text{electrical part}}$$

$\hat{v}_k$ : molar volume of species  $k$

$p_w$ : fluid (water) pressure in the pores

$R$ : universal gas constant

$T$ : absolute temperature

$x_k$ : the concentration of the species  $k$   
in the mixture

$z_k$ : is the valence of species  $k$

$F$ : Faraday's constant

$\phi$ : electrical field