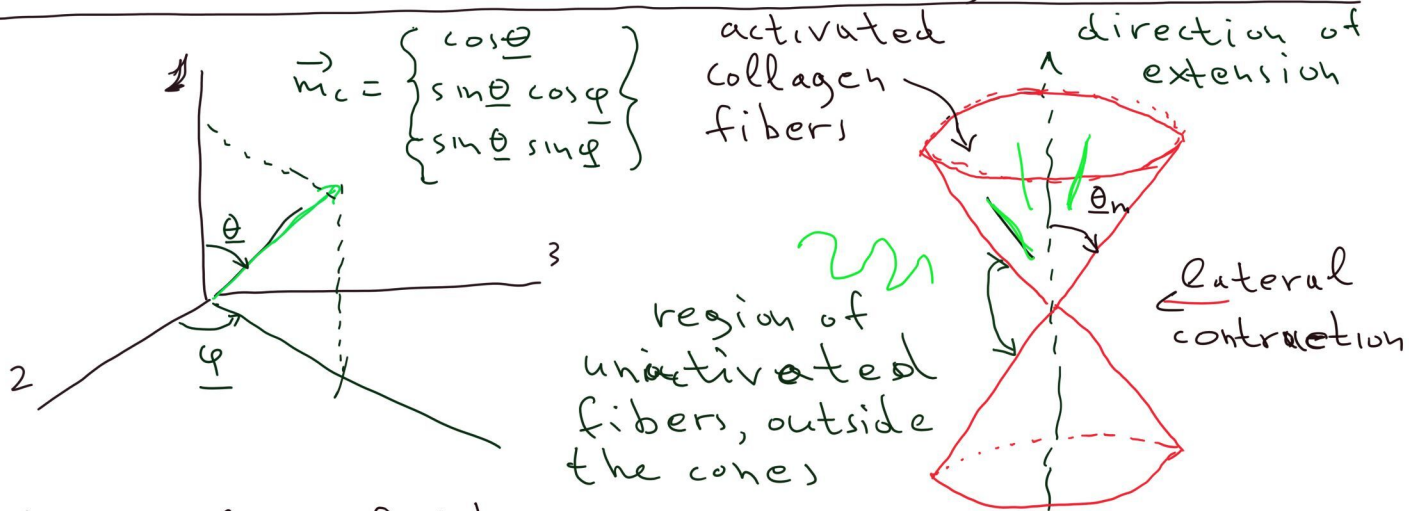


Random distribution of collagen fibers



θ_m : angle of the cone generator lines with the direction of extension. Fibers that lie within the cone, are activated.

For linear elastic behavior, we can have

$$w_c = \frac{K_c}{2} \langle \underline{\underline{\epsilon}} : \underline{\underline{M}}_c \rangle^2 \quad \left. \vphantom{\frac{K_c}{2}} \right\} \underline{\underline{\epsilon}}: \text{small strain tensor}$$

$$\frac{\partial w_c}{\partial \underline{\underline{\epsilon}}} = K_c \langle \underline{\underline{\epsilon}} : \underline{\underline{M}}_c \rangle \underline{\underline{M}}_c \quad \left\{ \begin{array}{l} \text{quantities for} \\ \text{a single fiber} \end{array} \right.$$

$$\frac{\partial^2 w_c}{\partial \underline{\underline{\epsilon}} \partial \underline{\underline{\epsilon}}} = K_c \underbrace{H(\underline{\underline{\epsilon}} : \underline{\underline{M}}_c)}_{\substack{\text{discontinuous} \\ \text{stiffness at zero} \\ \text{strain}}} \underline{\underline{M}}_c \otimes \underline{\underline{M}}_c$$

Suppose that all fibers are in extension. Integrate the quantities, referred to a single fiber, around the unit sphere and then take the average.

$$n^c \int_{S^2} w_c \frac{dS}{4\pi} = \frac{K^c}{15} \left[\frac{1}{2} (\underline{\underline{I}} : \underline{\underline{\epsilon}})^2 + \underline{\underline{\epsilon}} : \underline{\underline{\epsilon}} \right]$$

The integral $\times n^c$ gives the strain energy density, of the collagen part of the composite material, that consists of fibers embedded in a matrix.

n^c : volume density of the fibers in the composite material (fibers + matrix).

$$K^c = n^c K_c$$

Remember that for a uniaxial extension of a ~~bar~~ bar, consisting of Hookean linear material

$$w = \frac{1}{2} \sigma \varepsilon = \frac{1}{2} (E \varepsilon) \varepsilon = \frac{1}{2} E \varepsilon^2$$

Similarly, integrating around the unit sphere, for the other 2 quantities we get

$$n^c \int_{S^2} \frac{\partial w_c}{\partial \underline{\varepsilon}} \frac{dS}{4\pi} = \frac{K^c}{15} \left[(\underline{\varepsilon} : \underline{\underline{I}}) \underline{\underline{I}} + 2 \underline{\underline{\varepsilon}} \right] \rightarrow \text{stress}$$

$$n^c \int_{S^2} \frac{\partial^2 w_c}{\partial \underline{\varepsilon} \partial \underline{\varepsilon}} \frac{dS}{4\pi} = \frac{K^c}{15} \left[\underline{\underline{I}} \otimes \underline{\underline{I}} + 2 \underline{\underline{I}} \underline{\underline{\otimes}} \underline{\underline{I}} \right] \quad \underline{\underline{\otimes}} : \text{gives another 4th order tensor}$$

From the above relations, we can relate the Lamé constants, λ and μ , with the ^{single} fiber material constant K_c . We finally get

$$\lambda^c = \mu^c = \frac{K^c}{15}, \quad K^c = n^c K_c$$

Since $\lambda^c = \mu^c$, Poisson's ratio ~~can~~ takes the value 0,25.

If we have uniaxial extension along direction 1, and symmetry along direction 2 and 3, we get the activation relation

$$\underline{\varepsilon} : \underline{\underline{M}}_c \geq 0 \Rightarrow \underbrace{\cot \underline{\Theta} = \frac{m_{c1}^2}{m_{c2}^2 + m_{c3}^2}}_{\text{fiber orientation relation}} \geq \underbrace{\cot \underline{\Theta}_m = -\frac{\varepsilon_{22}}{\varepsilon_{11}}}_{\text{strain field relation}}$$

This is how $\underline{\Theta}_m$ can be evaluated.