

**NATIONAL TECHNICAL UNIVERSITY OF ATHENS** School of Civil Engineering – Geotechnical Department

# Computational Methods in the Analysis of Underground **Structures**

Spring Term 2023 – 24

Lecture Series in Postgraduate Programs:

- 1. Analysis and Design of Structures (DSAK)
- 2. Design and Construction of Underground Structures (SKYE)

Instructor: Michael Kavvadas, Emer. Professor NTUA

LECTURE 2: Stresses and deformations around a cylindrical tunnel (2D elasto-plastic analysis)

06.03.2024

#### Evolution of wall convergence along the tunnel axis (x)



#### NOTE: Floor rise is equal to crest settlement

- Convergence starts at distance 0.5-0.75 D ahead of the tunnel face
- 30% 50% of the total convergence has occurred at the tunnel face
- Wall convergence ceases to increase beyond about 1.5 D behind the tunnel face

Evolution of wall convergence along the tunnel axis (x)

Tunnel advance and wall support in steps with length (p).

The front part of the tunnel, close to the tunnel face (length  $d_1$ ), remains unsupported for construction purposes (access limitation of machinery). The maximum unsupported length close to the tunnel face is  $d_2 = d_1 + p$ 





Evolution of wall convergence along the tunnel axis

Wall convergence at the tunnel face (x=0) is about 31% of the maximum value

The maximum value increases in weaker ground, larger tunnel depth and larger tunnel size.

 $P_{o}$ 

### Evolution of wall convergence along the tunnel axis (x)



Empirical relationships btained from the esults of 3D finite lement analyses with wide range of round parameters nd tunnel geometries size and depth)

Evolution of wall convergence along the tunnel axis (x) (Chern, 1998)



$$
\frac{u_R(x)}{u_{R\infty}} = \left[1 + exp\left(0.91\frac{x}{R}\right)\right]^{-1.7}
$$

Chern (1998): Empirical relationship from the results of a set of 3D finite element analyses with a wide range of ground parameters and tunnel geometries (size and depth)

Evolution of wall convergence along the tunnel axis (x) (Chern, 1998)

Wall convergence  $u_R(x)$  of an unsupported tunnel at distance ( $x$ ) from the tunnel face (located at  $x = 0$ ):

$$
\frac{u_R(x)}{u_{R\infty}} = \left[1 + \exp\left(0.91\frac{x}{R}\right)\right]^{-1.7} \quad \text{or} \quad \frac{x}{R} = 1.10 \quad \ln\left[\frac{u_R(x)}{u_{R\infty}}\right]^{-0.588} - 1
$$

 $\mathsf{L}$ 

### *R* = tunnel radius

 $u_{R_{\infty}}$  = the final (maximum) wall convergence at large distance from the tunnel face  $(x = -\infty)$ . Can be calculated with analytical methods (present section), but more accurately with finite element analyses

 $u_R(o)$  = wall convergence at the tunnel face (location  $x$  = 0) According to Chern:  $u_R(o)$ =  $0.308$   $u_{R^{\infty}}$ 

 $\bullet$  2D model of tunnel excavation: The initial geostatic pressure  $(p_o)$  gradually reduces to (*p*) and eventually becomes zero. As the stress reduces, the tunnel wall converges (U<sub>R</sub>) up to a maximum value U<sub>R,max</sub> (when p=0). a cylindrical tunnel - Assumptions<br>nitial geostatic pressure  $(p_o)$  gradually<br>nes zero. As the stress reduces, the<br>ximum value  $U_{R,max}$  (when p=0).<br> $\lambda = 1 - \frac{p}{p_o} \implies p = (1 - \lambda) p_o$ <br>i ()<br>Wall convergence  $U_R$  reaches a<br>maximum Stresses and deformations around a cylindrical tunnel - Assumptions Deconfinement = Reduction of pressure p

Deconfinement coefficient:  $\lambda =$ 



$$
-\frac{p}{p_o} \Rightarrow p = (1 - \lambda) p_o
$$

lindrical tunnel - Assumptions<br>geostatic pressure  $(p_o)$  gradually<br>zero. As the stress reduces, the<br>im value  $U_{R,max}$  (when p=0).<br> $1 - \frac{p}{p_o} \Rightarrow p = (1 - \lambda) p_o$ <br>Wall convergence  $U_R$  reaches a<br>maximum value  $U_{R,max}$  and does not<br>c  $p$  - Assumptions<br> *p p p gradually*<br>
stress reduces, the<br>  $p = (1 - \lambda) p_o$ <br>  $p = (1 - \lambda) p_o$ <br>
nce  $U_R$  reaches a<br> *p* = ( $\theta_R$  reaches a<br> *p e D<sub>R,max</sub>* and does not<br>
crease more. Why ?<br>
tress change (*p* in 2D<br>
causes Wall convergence U<sub>R</sub> reaches a maximum value  $U_{R,max}$  and does not continue to increase more. Why ? Because the stress change (p in 2D models) that causes ground deformation, only occurs close to the tunnel face, i.e., along the length between  $\lambda=0$  and 1.

NOTE: Need a relation between p (i.e.,  $\lambda$ ) and location (x) to link 2D with 3D models

Relation  $U_R$  and p (or  $\lambda$ ): from 2D analysis (next) Relation  $U_R$  and x (from Chern)

Relation  $p$  (or  $\lambda$ ) and x

#### Stresses and deformations around a cylindrical tunnel - Assumptions

- 2D (plane) strain (no change along tunnel axis z)
- Cylindrical unsupported tunnel, with radius R
- Hydrostatic initial (geostatic) stress state (K<sub>o</sub> = 1  $\rightarrow \sigma_{\nu 0} = \sigma_{\hbar 0} = \rho_{\rho}$ )
- Elastic perfectly plastic ground, yielding with Mohr-Coulomb criterion (strength parameters:  $c, \varphi$ ) *volumetric strain*  $\delta \equiv \frac{\varepsilon_r + \varepsilon_\theta}{ } = \frac{$  $+\,\varepsilon_{\circ}\quad volumetr$
- Constant dilatancy (δ) in the plastic domain:  $tan \delta = \frac{r_1 + r_2}{r_1}$

$$
\equiv \frac{1}{\varepsilon_r - \varepsilon_\theta} = \frac{1}{\varepsilon_r}
$$



Stresses and deformations around a cylindrical tunnel - Assumptions

**Definitions:**

\nOverstress factor: 
$$
N_s = \frac{2p_o}{\sigma_{cm}}
$$
 (for  $K_o = 1$ )

\nGround strength:  $\sigma_{cm} = 2c\sqrt{k} = \frac{2c\cos\phi}{1-\sin\phi}$   $\overrightarrow{K_o p_o}$ 



Rockmass strength (empirical correlation with GSI):

 $\sum_{i=1}^n\frac{1}{i}$  $\left| \frac{\partial \omega}{\partial x_i} \right|$  $(25.5)$  $=\frac{\sigma_{ci}}{\sqrt{S}}$  exp $\left(\frac{GSI}{S}\right)$ 25.5  $\exp\left(-\frac{1}{2\pi\epsilon_0}\right)$  $50 \pm 25.5$  )  $G_{G}$   $GST$   $\sim$ *cm*  $\sigma_{\rm g}$  (g)  $\sim$  $\sigma_{\text{m}} = \frac{v}{\sqrt{2}} \exp(-\frac{1}{2}$ 





Stresses and deformations around a cylindrical tunnel - Assumptions

Ground deformations in plastic domain ( $r < r_p$ ): Dilatancy is constant in plastic domain (parameter K):

$$
\tan \delta = \frac{\varepsilon_r + \varepsilon_\theta}{\varepsilon_r - \varepsilon_\theta} \quad \text{where:} \quad K = \frac{1 + \tan \delta}{1 - \tan \delta} \quad \text{where}
$$

Strain definitions (u = radial displacement):

$$
\varepsilon_r = \frac{du}{dr} \qquad \varepsilon_0 = \frac{u}{r} \qquad \varepsilon_z = 0
$$



The above formulae give (inside the plastic domain (i.e., for  $r < r_{\rm p}$ ):

$$
\frac{du}{dr} + \frac{K}{r}u = 0 \Rightarrow u = cr^{-K} \Rightarrow c = u_p r_p^K \Rightarrow u = u_p \left(\frac{r_p}{r}\right)^K \text{ For } R < r < r_p
$$
\n
$$
(\text{at } r = r_p \Rightarrow u = u_p)
$$

 ${\sf u}_{\sf p}$  is the radial displacement at r =  ${\sf r}_{\sf p}^{\phantom{\dagger}}$  (calculated from the elastic zone)

At the tunnel wall (r = R) :  $u_p = u_p \left(\frac{r_p}{r}\right)^K$ *p*  $R \sim p \left(R\right)$ *r*  $u_R = u_p \left| \frac{P}{D} \right|$  $\int$  and  $\int$  $\bigwedge$  and  $\bigwedge$  $\begin{array}{|c|c|c|}\n\hline\nD\n\end{array}$  $(R)$  $\left(\frac{r_n}{r_n}\right)^K$ an di Kabupatén Ba

Stress-strain relationships in plane strain (cylindrical coordinates): bases and deformations around a cylindrical tunnel – Elastic domain (r > r<sub>p</sub>)<br> **K**<sub>c</sub> =  $\frac{1}{\Lambda} \left\{ \dot{c} \begin{pmatrix} \dot{c} \\ \dot{c} \end{pmatrix} \right\}$   $\Lambda = \frac{E}{(1+v)(1-v)}$   $K_o = \frac{v}{1-v}$ <br>  $\epsilon_e = \frac{1}{\Lambda} \left\{ \dot{\sigma}_e - K_o \dot{\sigma}_r \right\}$   $\dot{\sigma}_r = \sigma_r - p_o$ Stresses and deformations around a cylindrical tunnel – Elastic domain (r >  $\sf r_{\sf p}$ )



$$
\varepsilon_r = \frac{1}{\Lambda} \left\{ \dot{\mathbf{C}} \cdot \mathbf{K}_0 \right\} \qquad \Lambda \equiv \frac{E}{(1 + v)(1 - v)} \qquad K_o \equiv \frac{v}{1 - v}
$$
\n
$$
\varepsilon_{\theta} = \frac{1}{\Lambda} \left\{ \dot{\mathbf{C}}_{\theta} - K_o \dot{\mathbf{C}}_{r} \right\} \qquad \dot{\mathbf{C}}_{r} = \mathbf{C}_{r} - p_o \qquad \dot{\mathbf{C}}_{\theta} = \mathbf{C}_{\theta} - p_o
$$

 $\equiv$ 

 $(1-\nu)$ 

 $\frac{E(1-\nu)}{(1+\nu)(1-2\nu)}$ 

 $D \equiv \frac{E(1-\nu)}{2}$ 

 $d\dot{\sigma}_r$  ,  $\dot{\sigma}_r$  –  $\dot{\sigma}_\theta$  –  $\alpha$ 

*dr <sup>r</sup>*

 $1 + v(1 - 2v)$ 

 $+\frac{\sigma_r - \sigma_{\theta}}{1} = 0$ 

Solving for the stress increments:

$$
\dot{\sigma}_r = D\{\varepsilon_r + K_o \varepsilon_\theta\} \qquad \dot{\sigma}_\theta = D\{\varepsilon_\theta + K_o \varepsilon_r\} \qquad D = \frac{E(1-\varepsilon)}{(1+\nu)(1-\varepsilon)}
$$

Equilibrium equation (along axis r):

Strain definitions (u = radial displacement):

$$
\varepsilon_r = \frac{du}{dr} \qquad \varepsilon_\theta = \frac{u}{r} \qquad \varepsilon_z = 0 \qquad \Longrightarrow \qquad \frac{d^2u}{dr^2} + \frac{1}{r}\frac{du}{dr} - \frac{u}{r^2} = 0 \quad \Longrightarrow \qquad u = c_1r + \frac{c_2}{r}
$$

Boundary conditions:  $c_1 = 0$  (u cannot increase with r) If plastic zone exists: At  $r = r_p \rightarrow u = u_p \rightarrow c_2 = u_p r_p \rightarrow$ If plastic zone does not exist: At r=R  $\rightarrow$   $\sigma_{\text{r}}$ =p  $\rightarrow$ *p* | *f*: | | *p*<sub>1</sub>, 1</sub>,  $\mathbf{v}$ *r* 1 *u u r*  $(r_n)$  . . .  $\frac{1}{2} = u_p \left( \frac{p}{r} \right)$  (in elastic  $2^{-\kappa}$  2G  $\rightarrow$   $\cdots$  $c_0 = \lambda \frac{p_o R^2}{2}$   $\lambda u = \lambda R \frac{1}{2}$ *G* (2*G*)  $= \lambda \frac{P_0}{2G}$   $\rightarrow u = \lambda R \left( \frac{P_0}{2G} \right)^{-1}$  $u = \lambda R \left( \frac{P_o}{P} \right) \frac{R}{P}$ *G <sup>r</sup>*  $\lambda R \sim 1$  $= \lambda R \bigg(\frac{p_o}{2G}\bigg)\frac{R}{r}$ (in elastic zone)

Deformations around a cylindrical tunnel

The differential equation of equilibrium gives: *c c*<sub>o</sub>  $u = \frac{c_2}{c_1}$ 1. Linearly elastic ground,  $K_0 = 1$ 

Constant  $\mathbf{c}_2$  is determined from the stress boundary condition:  $\sigma_r(r=R) = p = p_o(1-\lambda)$ 

Thus, the radial displacement at distance (r) is:

Deformations around a cylindrical tunnel  
\nnearly elastic ground, K<sub>o</sub> = 1  
\ndifferential equation of equilibrium gives: 
$$
u = \frac{c_2}{r}
$$
  
\ntant c<sub>2</sub> is determined from the stress boundary condition:  
\n $\sigma_r (r = R) = p = p_o (1 - \lambda)$   
\nthe radial displacement at distance (r) is:  
\n $u = \lambda R \left(\frac{p_o}{2G}\right) \left(\frac{R}{r}\right) \implies u = \left(1 - \frac{p}{p_o}\right) R \left(\frac{p_o}{2G}\right) \left(\frac{R}{r}\right)$ NOTE:  
\ne tunnel wall (r=R):  $u_R = \lambda R \left(\frac{p_o}{2G}\right)$   
\nor complete deconfinement ( $\lambda = I, r=R$ ):  $u_{R\infty} = R \left(\frac{p_o}{2G}\right)$ 

At the tunnel wall  $(r=R):$   $\|u_R=\lambda\|R\| \frac{P_o}{\lambda}\|$  $\left( 2G\right)$  $=$   $\lambda$   $R\left(\frac{p_o}{p}\right)$   $\sim$ *p*<sub>0</sub> | *p*<sub>0</sub>  $u_{\nu} = \lambda R |\frac{F_{\nu}}{2}|$  $R - \frac{R}{2G}$  $\lambda$  R  $\frac{P_o}{\sim}$   $\sim$ 

and for complete deconfinement (  $\lambda=I,$   $r=R$  ) :  $\left| \begin{array}{c} u_{R \infty} = R \left| \begin{array}{c} P_{\text{o}} \end{array} \right| \right|$  $\infty$ <u> a se</u>

$$
\frac{u_R}{u_{R\infty}} = \lambda
$$
 Convel

Convergence-confinement curve in linearly elastic ground

=

 $\begin{picture}(20,20) \put(0,0){\dashbox{0.5}(5,0){ }} \put(15,0){\dashbox{0.5}(5,0){ }}$ ) and  $\overline{a}$ 

 $G$  = ground shear modulus  $R$  = tunnel radius ,  $p<sub>o</sub>$  = geostatic stress  $\lambda = \lambda(x)$  = deconfinement coefficient



NOTE: Strains and stresses calculated by differentiation

*r*

 $2(1+\nu)$ 

 $V$  ) and the set of the set of  $V$ 

 $+$   $\boldsymbol{\mathcal{V}}$  ) and the set of the set of  $\boldsymbol{\mathcal{V}}$ 



 $\lambda=1-$ 

 $0 < \lambda < 1$ 

*p<sup>o</sup>*

#### Stresses and deformations around a cylindrical tunnel – Only elasticity



**Linearity elastic ground**  
\n
$$
K_o = 1
$$
\n
$$
\sigma_r = p_o \left[ 1 - \lambda \left( \frac{R}{r} \right)^2 \right]
$$
\n
$$
\sigma_{\theta} = p_o \left[ 1 + \lambda \left( \frac{R}{r} \right)^2 \right]
$$
\n
$$
\lambda = 1 - \frac{p}{p_o}
$$

At tunnel wall (r=R):  $\sigma_r = p = (1 - \lambda)p_o$  $p = (1 - \lambda)p_o$  $\sigma_{\theta} = 2p_o - p = (1 + \lambda)p_o$ and for λ=1:  $\sigma_{\rm r} = 0$ ,  $\sigma_{\theta} = 2p_{\rm o}$ 

### Linearly elastic ground –  $K_0 \neq 1$  (Kirsch solution) Stresses and deformations around a cylindrical tunnel – Only elasticity

Circular tunnel (radius r<sub>o</sub>) at depth (H), unit weight of ground (γ), horizontal stress coefficient K ( $\sigma_h$  = K  $\sigma_v$ ). Geostatic stresses:  $\sigma_v$  =  $\gamma H$ ,  $\sigma_h$  = K  $\gamma H$  (do not vary with depth). Angle ( $\theta$ ) is measured from tunnel center, with respect to the vertical ( $\theta$ =0)

Tunnel is unsupported and  $\lambda = 1$  ( $\sigma_{rr} = 0$  at r =  $r_{o}$ )

#### **Kirsch solution** (for  $p=0$ ):

$$
\sigma_{rr} = \gamma H \left[ \frac{1+K}{2} \left( 1 - \frac{r_0^2}{r^2} \right) \right] + \gamma H \left[ \frac{1-K}{2} \left( 1 + 3\frac{r_0^4}{r^4} - 4\frac{r_0^2}{r^2} \right) \cos 2\vartheta \right]
$$
  

$$
\sigma_{\vartheta\vartheta} = \gamma H \left[ \frac{1+K}{2} \left( 1 + \frac{r_0^2}{r^2} \right) \right] - \gamma H \left[ \frac{1-K}{2} \left( 1 + 3\frac{r_0^4}{r^4} \right) \cos 2\vartheta \right]
$$
  

$$
\sigma_{r\vartheta} = -\gamma H \frac{1-K}{2} \left( 1 - 3\frac{r_0^4}{r^4} + 2\frac{r_0^2}{r^2} \right) \sin 2\vartheta
$$



Circumferential stress at springline (θ=90°):  $\sigma_{\theta\theta}$  = (3-K)γH  $-$  Initial value:  $\sigma_{\theta\theta}$  = γH For K =  $0.5 \rightarrow \sigma_{\theta\theta} = 2.5$  γH  $K = 1$  ->  $σ<sub>θθ</sub> = 2 yH$ 

Circumferential stress at crest and invert (θ=0 & 180°):  $\sigma_{\theta\theta}$  = (3K-1)γH - Initial value:  $\sigma_{\theta\theta}$  = KγH For K =  $0.5 \rightarrow \sigma_{\theta\theta} = 0.5 \text{ yH}$  (initial value 0.5 yH)  $K = 1$  ->  $σ_{\theta\theta} = 2$  γH (initial value γH)

# Linearly elastic ground –  $K_0 \neq 1$  (Kirsch solution) Stresses and deformations around a cylindrical tunnel – Only elasticity



Stresses and deformations around a cylindrical tunnel

2. Elastic – perfectly plastic ground,  $K_0 = 1$ 



The limit of the plastic zone (*rp*) depends on:

- •The tunnel radius (R)
- •the ground strength parameters (c,φ)
- the initial geostatic stress (*po*)
- the deconfinement coefficient (λ), i.e., the internal pressure (*p*)

# Convergence – confinement curve in elasto-plastic ground Influence of the σ-ε curve



The ground pressure (p) on the tunnel lining decreases with increasing tunnel wall convergence

# Convergence – confinement curve in elasto-plastic ground Influence of the σ-ε curve



If ground continuity is preserved, the convergence-confinement curve does NOT turn upward (collapse) even in strongly strain softening ground. If, however, ground continuity is lost (e.g. rock block contact is lost) due to large ground deformations, then the convergence-confinement curve may turn upwards (collapse).

This means that ground pressure on the tunnel lining will increase at large ground deformations.

#### Stresses around a cylindrical tunnel

2. Elastic – perfectly plastic ground,  $K_0 = 1$ 

Calculation of the minimum internal pressure  $p = p_{cr}$  to maintain elasticity in the ground:

Stress distribution in the elastic domain:

$$
\sigma_r = p_o \left[ 1 - \lambda \left( \frac{R}{r} \right)^2 \right] \qquad \sigma_\theta = p_o \left[ 1 + \lambda \left( \frac{R}{r} \right)^2 \right] \qquad \lambda = 1 - \frac{p_{cr}}{p_o}
$$

Stresses (elastic) at the tunnel wall (r=R):

 $\sigma_3 = \sigma_r = p_{cr}$ Marginal fulfillment of the M-C failure criterion at the tunnel wall:

$$
\sigma_1 = k \sigma_3 + \sigma_{cm} \quad \implies \quad \frac{p_{cr}}{p_o} = \left(\frac{2}{1+k}\right) \left(\frac{N_s - 1}{N_s}\right)
$$

Critical deconfinement coefficient:

CONCLUSION: There is no plastic zone around the tunnel, If:  $\Lambda_{cr} \ge 1$  (i.e.,  $N_s \le 1$ ) or if:  $\lambda_{cr}$  < 1 and  $\lambda \leq \lambda_{cr}$ 

Plastic zone develops around the tunnel if:  $\lambda_{cr}$  < 1 (i.e., N<sub>s</sub> > 1) and  $\lambda > \lambda_{cr}$ 

$$
\begin{array}{ccc}\n & \text{ELASTIC} & \text{PLASTIC} \\
\lambda=0 & \lambda_{cr} & \lambda=1\n\end{array}
$$



**Contract Contract** и производство на селото на се<br>Становите селото на селото на

 $\frac{1}{N}$  $\left(\begin{array}{cc} N_s \end{array}\right)$ 

*s*

*s* / *s* 

 $\begin{array}{|c|c|c|c|c|c|}\n\hline\n & & S & 1 \\
\hline\n & & 1 & 1\n\end{array}$  $\left\langle N_{s}\right\rangle$ 

 $\bigvee N_{\epsilon}-1$ 

 $+$  K  $\land$  N  $\rightarrow$  N  $\rightarrow$ 

 $p_{cr}$  1  $\left(2 \right) \left\langle N_s - 1 \right\rangle$ 

 $\sum_{i=1}^n a_i$ 

 $(N, -1)$ 

*N*

 $\sigma_1 = \sigma_\theta = 2 p_o - p_{cr}$ 

 $\left| \frac{Z}{Z} \right| \frac{1 \sqrt{S}}{Z}$  $(1+k)$   $N_s$ 

 $p_{\scriptscriptstyle o}$   $(1+k)$   $N_{\scriptscriptstyle s}$   $)$ 

*cr D*  $\left(1+k \sqrt{N}\right)$   $\left(1+2k \sqrt{N}\right)$ 

 $1+k$   $N$   $\quad$   $\quad$ 

 $\lambda_{-} = 1 - \frac{p_{cr}}{r} = 1 - \left(\frac{2}{r_{-}}\right)\left(\frac{N_s - 1}{r_s}\right)$ 

*o*

*cr*

 $\left( 2 \right) \left( N_{s} -$ 

#### Stresses around a cylindrical tunnel – elastoplastic ground

Critical deconfinement coefficient – ground remains elastic but M-C failure criterion is marginally fulfilled at the tunnel wall (i.e.,  $r_p = R$ ):

$$
\left| \lambda_{cr} = 1 - \frac{p_{cr}}{p_o} \implies \lambda_{cr} = 1 - \left( \frac{2}{1+k} \right) \left( \frac{N_s - 1}{N_s} \right) \right| \qquad \lambda_s = \tan^2 \left( 45 + \frac{\varphi}{2} \right)
$$

$$
N_s = 2p_o / \sigma_{cm}
$$
  

$$
k = \tan^2 \left( 45 + \frac{\varphi}{2} \right)
$$
  

$$
\sigma_{cm} = \frac{\sigma_{ci}}{52.63} \exp \left( \frac{GSI}{20} \right)
$$

Values of  $\lambda_{cr}$  (plastic zone around the tunnel develops if  $\lambda > \lambda_{cr}$ )



Stresses around a cylindrical tunnel – elastoplastic ground

Example: γ = 22 kN/m<sup>3</sup>, H = 100 m, K<sub>o</sub> = 0.60 ⇒ p<sub>o</sub> = 0.5 (1+K<sub>o</sub>) γ H = 1.76 MPa GSI = 25,  $\sigma_{ci}$  = 12 MPa, E<sub>i</sub> = 13.5 GPa  $\Rightarrow \sigma_{cm}$  = 0.64 MPa, E = 821 MPa  $v = 0.30 \Rightarrow G = 316 \text{ MPa}$  $\overline{\varphi} = 32^{\circ} \Rightarrow k = 3.2546$  $\delta$  = 7°  $\Rightarrow$  K = 1.28  $\exp\left(-\frac{1}{2} \right)$ 52.63  $\sqrt{20}$ *ci*  $cm \sim$   $\sim$   $\sim$  $\sigma$  (*GSI*)  $\sigma_{\dots} = \frac{c}{c}$  exp  $rac{E}{2(1+v)}$   $\sigma_{cm} = \frac{\sigma_{ci}}{52.63} \exp\left(\frac{GSI}{20}\right)$  $D =$  damage factor  $(=0)$  $G = \frac{E}{\sqrt{E}}$   $G = \frac{C}{\sqrt{E}}$  $V$   $\sim$   $\sim$   $\sim$   $\sim$   $\sim$  $=\frac{1}{2(1+\nu)}$   $\sigma_{cm} = \frac{1}{5}$  $K \equiv \frac{1}{2}$  $\sqrt{1-t}$  $\sum_{\mathbf{r}}$  1+ta  $\left| 45 + \frac{\varphi}{2} \right|$  K  $\binom{1}{1}$  2  $=$  tan<sup>2</sup> $\left(45+\frac{\varphi}{\epsilon}\right)$   $K \equiv \frac{1+\tan \pi}{\epsilon}$ 2  $l$   $l$   $-$  tan $l$  $k = \tan^2 \left( 45 + \frac{\varphi}{\pi} \right)$   $K = \frac{1 + \tan \theta}{\pi}$ 1 – tan $\delta$  and the set of  $\delta$  $E_{\rm rm} = E_{\rm i} \left( 0.02 + \frac{1 - D/2}{1 + e^{((60 + 15D - \text{GSD})/11)}} \right)$  $1+ \tan\delta$  and the state of  $\delta$  $K = \frac{1 + \tan \phi}{\sqrt{1 - \frac{1}{\phi^2}}}$ 

Calculations:

$$
N_s = \frac{2 p_o}{\sigma_{cm}} = 5.5 \qquad \lambda_{cr} = 1 - \left(\frac{2}{1+k}\right) \left(\frac{N_s - 1}{N_s}\right) = 0.615
$$

Result: For  $\lambda > 0.615$ , i.e., for  $p/p_0 < 0.385$  plastic zone develops around the tunnel



# CASE 1: Ground remains elastic (no plastic zone)  $\lambda_{cr} = 1 - \left(\frac{2}{1+k}\right)\left(\frac{N_s-1}{N_s}\right)$

- If  $N_s \leq 1$   $\rightarrow$  for all  $\lambda$
- If  $N_s > 1 \rightarrow$  for  $\lambda \leq \lambda_{cr}$

Stresses around the tunnel:

$$
\sigma_r = p_o \left[ 1 - \lambda \left( \frac{R}{r} \right)^2 \right]
$$

$$
\sigma_\theta = p_o \left[ 1 + \lambda \left( \frac{R}{r} \right)^2 \right]
$$

Displacement around the tunnel:

$$
u = \lambda \ R \left(\frac{p_o}{2G}\right) \left(\frac{R}{r}\right) \qquad \qquad \sum_{n=1}^{\infty} \sum_{i=1}^{\infty} \frac{p_i}{n}
$$

At the tunnel wall (*r=R*):

$$
u_R = \lambda \ R \left( \frac{p_o}{2G} \right)
$$





CASE 2: Plastic zone develops in the ground  
\nIf N<sub>s</sub> > 1 and 
$$
\lambda > \lambda_{cr}
$$
  
\n
$$
\frac{EIASTIC}{\lambda_{c0}} = \frac{0}{R} = \frac{PIASTIC}{\lambda_{cr}}
$$
\nRadius of plastic zone  $(r_p)$ :  
\n1. If  $\varphi = 0$ :  $\frac{r_p}{R} = \exp\left[\frac{1}{2}(\lambda N_s - 1)\right]$   
\n2. If  $\varphi > 0$ :  $\frac{r_p}{R} = \left\{\left(\frac{2}{k+1}\right)\left[\frac{2+N_s(k-1)}{2+N_s(k-1)(1-\lambda)}\right]\right\}^{\frac{1}{k-1}}$   
\nAnd in full deconfinement  $(\lambda = 1)$ :  $\frac{r_{p\infty}}{R} = \left\{\left(\frac{1}{k+1}\right)[2+N_s(k-1)]\right\}^{\frac{1}{k-1}}$ 

# Proof of formulae for  $\sigma_r$  and  $\sigma_\theta$ :

Επίλυση στην πλαστική ζώνη, δηλαδή για R < r < r<sub>p</sub>

Εξίσωση ισορροπίας:

$$
\frac{d\sigma_r}{dr} + \frac{\sigma_r - \sigma_\theta}{r} = 0
$$

Κριτήριο αστοχίας Mohr-Coulomb :  $\sigma_{\theta} = k \sigma_r + \sigma_{cm}$ 

Απαλειφή του σ<sub>θ</sub> δίνει:  $\frac{a}{2}$ 

$$
\frac{d\sigma_r}{dr} - \frac{1}{r}(k-1)\sigma_r - \frac{1}{r}\sigma_{cm} = 0
$$

Επίλυση της ανωτέρω διαφορικής εξίσωσης:

(α) Περίπτωση:  $k \neq 1 \Rightarrow \varphi \neq 0$ :

Με συνοριακή συνθήκη:  $\sigma_r(r=R) = p = (1 - \lambda)p_o$ 

$$
\sigma_r = \left[ (1 - \lambda)p_o + \left(\frac{\sigma_{cm}}{k - 1}\right) \right] \left(\frac{r}{R}\right)^{k - 1} - \left(\frac{\sigma_{cm}}{k - 1}\right)
$$

$$
\delta \eta \lambda \alpha \delta \eta : \frac{\sigma_r}{p_o} = \left[ (1 - \lambda) + \frac{2}{(k - 1)N_s} \right] \left(\frac{r}{R}\right)^{k - 1} - \frac{2}{(k - 1)N_s}
$$

$$
\sigma_\theta = k \sigma_r + \sigma_{cm} \implies \frac{\sigma_\theta}{p_o} = k \frac{\sigma_r}{p_o} + \frac{2}{N_s}
$$





# Proof of formulae for  $r_p/R$ :

Εξίσωση των τιμών των σ<sub>ε</sub> και σ<sub>θ</sub> στο όριο μεταξύ ελαστικής και πλαστικής ζώνης (r =  $r_p$ ) δίνει τις τιμές των  $c_2$  και  $r_p$ :

(α) Περίπτωση  $k \neq 1 \Rightarrow \varphi \neq 0$ :

$$
\frac{\sigma_r}{p_o} = \left[ (1 - \lambda) + \frac{2}{(k - 1)N_s} \right] \left( \frac{r_p}{R} \right)^{k - 1} - \frac{2}{(k - 1)N_s} = 1 - c_2 \left( \frac{2G}{p_o} \right) \frac{1}{r_p^2}
$$

$$
\frac{\sigma_o}{p_o} = k \frac{\sigma_r}{p_o} + \frac{2}{N_s} = 1 + c_2 \left( \frac{2G}{p_o} \right) \frac{1}{r_p^2}
$$

$$
\text{orote: } \frac{r_p}{R} = \left[ \left( \frac{2}{k + 1} \right) \frac{N_s + \frac{2}{k - 1}}{(1 - \lambda)N_s + \frac{2}{k - 1}} \right]^{\frac{1}{k - 1}}
$$

(β) Περίπτωση :  $k = 1 \Rightarrow \varphi = 0$  :







Stresses around a cylindrical tunnel – elastoplastic ground

# Ground remains elastic, always (for all  $\lambda$ ) if:  $\;\;N_{_S} \leq 1$

Schematic size of the plastic zone  $(r_p)$  around the tunnel



### Stresses around a cylindrical tunnel – elastoplastic ground Classification of tunnel excavation problems with  $N_{s}$  value



Hoek E and Marinos P 2000 Predicting Tunnel Squeezing Problems in Weak Heterogeneous Rock Masses. Tunnels and Tunnelling International 32(11) 45-51

#### Stresses around a cylindrical tunnel – elastoplastic ground

#### Radius of the plastic zone  $r_p$  (unsupported tunnel)



*p<sup>o</sup> =* geostatic stress *λ* = deconfinement coefficient *σcm* = ground strength *Ν<sup>s</sup>* = overstress factor

$$
p = p_o(1 - \lambda)
$$

$$
N_s = \frac{2 p_o}{\sigma_{cm}}
$$

*Graph is valid for common values of the relevant parameters*

Stresses around a cylindrical tunnel – elastoplastic ground

# Radius of the plastic zone  $r_p$  (supported tunnel,  $p_i$  = support pressure)



 $p_o$  = geostatic stress *λ* = deconfinement coefficient *σcm* = ground strength *Ν<sup>s</sup>* = overstress factor

$$
p = p_o(1 - \lambda)
$$



*Graph is valid for common values of the relevant parameters*

# CASE 2: Plastic zone develops in the ground

If N<sub>s</sub> > 1 and  $\lambda$  >  $\lambda_{\text{cr}}$ 

**ELASTIC PLASTIC**  $\lambda = 0$  $\lambda = 1$  $\lambda_{\rm cr}$ 



<mark>a</mark> sa salawan и производство и село в с<br>Село в село в село

<mark>a</mark> sa salawan и производство и село в с<br>Село в село в село

*r*

*R*  $\lfloor r \rfloor$  . The set of  $\lfloor r \rfloor$ 

 $\begin{array}{|c|c|c|c|c|}\n\hline\n\textbf{r} & \textbf{r} & \textbf{r} & \textbf{r} \\
\hline\n\textbf{r} & \textbf{r} & \textbf{r} & \textbf{r} & \textbf{r} \\
\hline\n\textbf{r} & \textbf{r} & \textbf{r} & \textbf{r} & \textbf{r} & \textbf{r}\n\end{array}$  $\binom{r}{r}$ 

 $\left(\frac{r_n}{r_n}\right)^{\Lambda}$ 

*r*

*K*

**Contract Contract June 1999** 

 $\bigwedge$  and  $\bigwedge$  and  $\bigwedge$  and  $\bigwedge$ 

 $\binom{r}{r}$ 

*u v v v* 

*p* | *p* |

*R*

 $u_n$   $u_n$ 

 $\overline{\phantom{a}}$ 

*r*

*r*

 $\binom{r_n}{r_n}$ 

 $\sum_{i=1}^n a_i$ 

 $\sum_{i=1}^n a_i$ 

Ground displacement :

(a) Displacement ( $u_p$ ) at the limit of the plastic zone (r =  $r_p$ ) :

Calculated for a tunnel with radius  $R=r_p$  and critical deconfinement ( $\lambda_{cr}$ ), in which case ground displacements are elastic for r >  $\sf r_{\sf p}$  :

$$
\frac{u_p}{R} = \lambda_{cr} \left(\frac{r_p}{R}\right) \left(\frac{p_o}{2G}\right) \qquad \lambda_{cr} = 1 - \left(\frac{2}{1+k}\right) \left(\frac{N_s - 1}{N_s}\right)
$$

(b) Displacement (u) in the elastic zone (r > r<sub>p</sub>) :  $u = u_{p} \left| \frac{p}{n} \right|$ Calculated by the elastic formula: <u>a matana</u>  $u = u$   $\frac{p}{q}$ 

*r c*  $u = \frac{c_2}{c_1}$  with boundary condition:  $u = u_p$  at  $r=r_p$ (c) Displacement (u) in the plastic zone  $(r < r_p)$ :  $\frac{a}{p} = \frac{p}{p} \left| \frac{p}{p} \right|$ and at the tunnel wall ( $r = R$ ) :  $\frac{R}{P} = \frac{P}{P}$   $\frac{p}{P}$ *K R r R u R*  $u_n$ ,  $u_n$ ,  $r_n$ ,  $r_n$ **Contract Contract** и производство и село в с<br>Село в село  $\begin{array}{|c|c|c|c|c|}\n\hline\nD & \mbox{ } & \mbox{ } \n\end{array}$  $(R)$  $\left(r_{n}\right)^{\Lambda}$  $\equiv$ 

(d) Displacement (u<sub>R∞</sub>) at tunnel wall at full deconfinement (λ=1):

$$
\frac{u_{R\infty}}{R} = \lambda_{cr} \left(\frac{p_o}{2G}\right) \left(\frac{r_{p\infty}}{R}\right)^{K+1} \lambda_{cr} = 1 - \left(\frac{2}{1+k}\right) \left(\frac{N_s - 1}{N_s}\right)
$$
\n
$$
\frac{r_{p\infty}}{R} = \left\{ \left(\frac{1}{k+1}\right) \left[2 + N_s(k-1)\right] \right\}^{\frac{1}{k-1}}
$$
\ni.e.,: 
$$
\frac{u_{R\infty}}{R} = f\left(\frac{p_o}{2G}, N_s, \phi, \delta\right)
$$

Displacement (u<sub>R</sub>) at tunnel wall, for any deconfinement  $\lambda > \lambda_{cr}$ ):

$$
\frac{u_R}{u_{R\infty}} = \left\{\frac{1}{1 + \frac{N_s}{2}(k-1)(1-\lambda)}\right\}^{\left(\frac{K+1}{k-1}\right)} = f(\lambda)
$$

$$
\begin{cases}\n\left(\frac{K+1}{k-1}\right) \\
= f\left(\lambda \right), N_s, \phi, \delta\n\end{cases}
$$

Proof of formulae for  $u_r$  :

**Β.3 Υπολογισμός των μετακινήσεων στην πλαστική ζώνη (r < rp ) :** 

Ορισμός διαστολικότητας στην πλαστική ζώνη :  $\tan \delta = \frac{\varepsilon_r + \varepsilon_\theta}{\varepsilon_r - \varepsilon_a} \ge 0$ 

$$
\text{otherwise:} \quad K \equiv \frac{1 + \tan \delta}{1 - \tan \delta} = -\frac{\varepsilon_r}{\varepsilon_\theta} \ge
$$

$$
A\lambda\lambda\dot{\alpha}:\quad \varepsilon_r=\frac{du}{dr}\quad,\ \varepsilon_\theta=\frac{u}{r}
$$

Orπότε: 
$$
\varepsilon_{\theta} K + \varepsilon_r = 0
$$
 ⇒  $\frac{u}{r} K + \frac{du}{dr} = 0$  ⇒  $u = \alpha \frac{1}{r} K$ 

Συνοριακή συνθήκη:  $r = r_p \implies u = u_p \implies u = u_p \left(\frac{r_p}{r}\right)^2$ 

Αλλά το  $u_p$  έχει υπολογισθεί από την ελαστική ζώνη. Συνεπώς:

(α) Περίπτωση  $k \neq 1 \Rightarrow \varphi \neq 0$ :

$$
u = u_p \left(\frac{r_p}{r}\right)^K \Rightarrow \frac{u_R}{R} = \frac{u_p}{R} \left(\frac{r_p}{R}\right)^K
$$

όπου:

$$
u_p = r_p \left(\frac{p_o}{2G}\right)\left(1 - \frac{2}{k+1}\right)\left[1 + \frac{2}{(k-1)N_s}\right] \implies u_p = r_p \left(\frac{p_o}{2G}\right)\frac{(k-1)N_s + 2}{(k+1)N_s}
$$

Proof of formulae for  $u_r$ :

$$
\kappa \alpha l : \frac{r_p}{R} = \left[ \left( \frac{2}{k+1} \right) \frac{N_s + \frac{2}{k-1}}{(1-\lambda)N_s + \frac{2}{k-1}} \right]^{k-1}
$$
  
  

$$
\Gamma \alpha \lambda = 1 :
$$
  

$$
\frac{r_{p\infty}}{R} = \left[ \frac{(k-1)N_s + 2}{k+1} \right]^{k-1} \kappa \alpha l \frac{u_{p\infty}}{R} = \frac{r_{p\infty}}{R} \left( \frac{p_o}{2G} \right) \frac{(k-1)N_s + 2}{(k+1)N_s}
$$

 $-1$ 

Προσδιορισμός της τελικής (για λ=1) σύγκλισης του τοιχώματος της σήραγγας  $(u_{\kappa_{\infty}})$  :

$$
\frac{u_{R\infty}}{R} = \frac{u_{P\infty}}{R} \left(\frac{r_{P\infty}}{R}\right)^K \Rightarrow \frac{u_{R\infty}}{R} = \frac{1}{N_s} \left(\frac{p_o}{2G}\right) \left[\frac{(k-1)N_s + 2}{k+1}\right]^{\frac{K+k}{k-1}}
$$

Παρατήρηση:

Επειδή στην ελαστική περίπτωση η τελική (για λ=1) σύγκλιση του

τοιχώματος της σήραγγας ( $u_{R\infty,e}$ ) είναι :  $\displaystyle{\frac{u_{R\infty,e}}{R} \!=\!\! \left(\frac{p_o}{2G}\right)}$  προκύπτει ότι :

$$
\frac{u_{R\infty}}{u_{R\infty,e}} = \frac{1}{N_s} \left[ \frac{(k-1)N_s + 2}{k+1} \right]^{\frac{K+k}{k-1}} \quad \text{if} \quad \frac{u_{R\infty}}{u_{R\infty,e}} = \frac{1}{N_s} \left[ \frac{r_{p\infty}}{R} \right]^{K+k}
$$

Proof of formulae for u<sub>r</sub>:

$$
u = u_p \left(\frac{r_p}{r}\right)^K \Rightarrow \frac{u_R}{R} = \frac{u_p}{R} \left(\frac{r_p}{R}\right)^K
$$
  
όπου:  $u_p = r_p \left(\frac{p_o}{2G}\right) \frac{1}{N_s}$   
\n
$$
\kappa \alpha \text{I:} \quad r_p = R \, \exp\left[\frac{1}{2}(N_s \lambda - 1)\right] \kappa \alpha \text{I} \quad r_{p\infty} = R \, \exp\left[\frac{1}{2}(N_s - 1)\right]
$$
\n  
\nΠροσδιορισμός της τελικής (για λ=1) σύγκλισης του τοιχώματος της  
\nσήραγγας (*u<sub>R</sub>*<sub>0</sub>) :

$$
\frac{u_{R\infty}}{R} = \frac{u_{p\infty}}{R} \left(\frac{r_{p\infty}}{R}\right)^K \Rightarrow
$$
\n
$$
\frac{u_{R\infty}}{R} = \frac{1}{N_s} \left(\frac{p_o}{2G}\right) \exp\left[\frac{1}{2}(N_s - 1)(K + 1)\right]
$$

$$
\delta \eta \lambda \alpha \delta \eta : \frac{u_{R\infty}}{u_{R\infty,e}} = \frac{1}{N_s} \exp \left[ \frac{1}{2} (N_s - 1)(K + 1) \right] \quad \eta \frac{u_{R\infty}}{u_{R\infty,e}} = \frac{1}{N_s} \left[ \frac{r_{p\infty}}{R} \right]
$$

 $\neg K+1$ 

# Displacement of tunnel wall *(u<sup>R</sup> )*

#### *Unsupported tunnel*

#### *Supported tunnel*



• *Data points for common values of the relevant parameters*

• *Significant reduction of wall displacement with increasing support pressure*

# Stresses at the tunnel wall (r=R)



Combining the Chern  $u_R$ -x curve (1), with the convergence-confinement  $u_R$ -p curve (2), one can develop the Chern-Panet curves λ-x (3). The Chern-Panet curves are useful in 3D numerical analyses (to compute λ from x)



### Chern-Panet curves

1. Displacement  $u<sub>R</sub>$  at tunnel wall as a function of deconfinement ( $\lambda$ ):

 $\left|\frac{K+1}{2}\right|$  $\left(k-1\right)$  $(K+1)$ 

For deconfinement  $\lambda > \lambda_{cr}$  (plasticity):

 $2 \left[ \begin{array}{ccc} 2 & 2 & 2 \end{array} \right]$ 

 $+$   $(k - 1)$   $\lambda$  )

 $1+\frac{1+s}{s}(k-1)(1-1)$ 

 $1+\frac{N_s}{(k-1)(1-1)}$  $\begin{pmatrix} 1 & 2 & \sqrt{k} & 1 \end{pmatrix}$ 

二 く<del>ーーーー、</del> ――――――

 $\sim$   $\sim$   $\sim$   $\sim$   $\sim$   $\sim$  $\left\{\frac{1}{N}\right\}$ 

 $\int_{\mathbb{R}^d} |u(x)|^2 dx$ 

 $(\lambda)$  and  $\lambda$  and  $\lambda$ 

 $u_{\rm n}(\lambda)$  and  $u_{\rm n}(\lambda)$ 

 $\lambda$  and  $\lambda$  and  $\lambda$ 

For deconfinement  $\lambda < \lambda_{cr}$ (elasticity):

**Contract Contract** 

$$
\frac{u_R(\lambda)}{u_{R\infty}} = \frac{1}{1 + \frac{N_s}{2}(k-1)(1-\lambda)} \qquad \qquad = f(\lambda \ ; N_s \ , \phi \ , \delta \ ) \qquad \qquad \frac{u_R}{u}
$$

$$
\frac{u_R(\lambda)}{u_{R_{\infty}}} = \lambda
$$

2. Displacement  $u_R$  at tunnel wall along the tunnel axis (x) (Chern, 1998) :  $\left(x\right)$   $\left[\begin{array}{c} \overline{1} & \overline{1} & \overline{1} \\ 1 & \overline{1} & \overline{1} \end{array}\right]$ 1.7  $1 + exp|0.91 - 1$  $\infty$ **Contract Contract** a sa salawan  $\int_{0}^{\frac{1}{2}}$  $1 + exp 0.9$  $\begin{array}{ccc} \begin{array}{ccc} \end{array} & \begin{array}{c} \end{array} & \begin{array}{c} \end{array} \end{array}$  $\begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$ **Maria Bara**   $\bigcap$   $\bigcap$  $|0.91 \frac{\lambda}{2}|$  $\begin{pmatrix} 1 & R & R & R \end{pmatrix}$  $\left( \begin{array}{c} x \\ y \end{array} \right)$  $= 1 + exp[0.9 - 1]$ *R x* 1 and *x* 1 and *x* 1 and *x* 1 and *x exp* |  $\cup$ .91  $\pm$  | |  $u_{\rm max}$  . The set of  $u_{\rm max}$  $u_n(x)$  |  $\ldots$  | *R∞* Letter in  $R \vee V$  1 1

je poznata za ostali se ostali<br>Dogodki se ostali se

1 *k* 1 ) and the second state  $\sim$ 

 $\left(\overline{k-1}\right)$ 

 $K+1$ ) and  $K+1$ 

 $+$  100  $-$  100  $-$  100  $-$ 

**Contract Contract** Jan Barbara (1980)  $\sum_{i=1}^n a_i$ 

Combination of (1), (2) gives the Chern-Panet curves, in the form:

$$
\lambda = f\left(\frac{x}{R}; N_s, \varphi, \delta\right)
$$

These curves calculate the deconfinement coefficient  $(\lambda)$  at any location  $(x)$  along the tunnel axis. They are used in numerical analyses for the calculation of (λ) at the location (x) of support application

Example:

 $\gamma$  = 22 kN/m $^3$ , K $_{\rm o}$  = 0.60  $\implies$  p $_{\rm o}$  = 0.5 (1+K $_{\rm o}$ )  $\gamma$  H = 1.76 MPa  $v = 0.30 \Rightarrow G = 316 \text{ MPa}$   $G = \frac{E}{2(1 + 1)}$  $\overline{\varphi} = 32^{\circ} \Rightarrow k = 3.2546$  $\delta$  = 7°  $\Rightarrow$  K = 1.28 Tunnel radius  $D = 6m -$  tunnel depth  $H = 100m$ <u>in de la provincia de la p</u>  $\begin{array}{c} +\frac{1}{2} \\ \end{array}$   $45 + \frac{\varphi}{2}$  $\binom{16}{1}$  2  $k = \tan^2\left(45 + \frac{\varphi}{\sqrt{2}}\right)$ 2 decree  $\sim$   $\sim$  $\delta$  and  $\delta$  in the set of  $\delta$  $1 + \tan \delta$  and  $\sin \theta$  $K \equiv \frac{1 + \tan \phi}{\sqrt{1 - \frac{1}{\phi^2}}}$ GSI = 25,  $\sigma_{ci}$  = 12 MPa, E<sub>i</sub> = 13.5 GPa  $\Rightarrow \sigma_{cm}$  = 0.64 MPa, E = 821 MPa  $\exp\left(-\frac{1}{2} \right)$ 52.63  $\rightarrow$  20  $\rightarrow$ *ci*  $cm \sim$   $\sim$   $\sim$   $\sim$   $\sim$  $\sigma_{\rm s}$  *(GSI*)  $\sigma_{\mu} = \frac{a}{2}$  expi - $\frac{E}{2(1+\nu)}$   $\sigma_{cm} = \frac{\sigma_{ci}}{52.63} \exp\left(\frac{GSI}{20}\right)$ ╊

> $\delta$  and  $\delta$  and  $\delta$  $1 - \tan \delta$  and  $\sim$  1  $D =$  damage factor  $(=0)$

Calculations:

$$
N_s = \frac{2 p_o}{\sigma_{cm}} = 5.5 \qquad \lambda_{cr} = 1 - \left(\frac{2}{1+k}\right) \left(\frac{N_s - 1}{N_s}\right) = 0.615
$$

$$
\frac{r_{p\infty}}{R} = \left\{ \left( \frac{1}{k+1} \right) \left[ 2 + N_s \left( k-1 \right) \right] \right\}^{\frac{1}{k-1}} = 1.72 \qquad \frac{u_{R\infty}}{R} = \lambda_{cr} \left( \frac{p_o}{2G} \right) \left( \frac{r_{p\infty}}{R} \right)^{K+1}
$$

 $u_{\text{max}} = 0.000 \times 0.000 \times 0.000 \times 0.000$ *R*  $u_{\rm m}$  , we have a set  $u_{\rm m}$  $R\infty$   $\cup$   $\cup$   $\sim$   $\sim$  $R^{\infty}$   $\Omega^{\infty}$  $=0.00588 \, \Rightarrow \, u_{{}_{R\infty}}=600\!\times\!0.00588\!=\!3.53$  $\frac{\infty}{6}$  = 0.00588  $\Rightarrow$   $u_{n_{1}}$  = 600 × 0.00588 = 3.53 cm

 $GRC =$  Ground Reaction Curve:  $\bigcup_{r}(x) / U_{r,max}$  versus  $p_i / p_o$ LDP = Longitudinal Displacement Profile or Chern curve: U<sub>r</sub>(x) / U<sub>r,max</sub> versus x / R Combination of the GRC and LDP curves provides the relation:  $p_i / p_o$  versus  $x / R$ which is required in 2D numerical analyses. Convergence - Confinement (or Ground Reaction) curve: Ground Reaction and Longitudinal Displacement (Chern) curves



EXCEL spreadsheet for the calculation of the GRC and LDP curves

Input data: R, p<sub>o</sub>, 
$$
\sigma_{cm}
$$
,  $\varphi$ ,  $\delta$ , G  
\nCalculate N<sub>s</sub>, k, K and  $\lambda_{cr}$   
\n
$$
K = \frac{1 + \tan \delta}{1 - \tan \delta} \quad \lambda_{cr} = 1 - \left(\frac{2}{1 + k}\right) \left(\frac{N_s - 1}{N_s}\right)
$$
\n
$$
Col 1: p/po between 1 ... 0\nCol 2:  $\lambda$  (between 0 ... 1)  
\n
$$
\lambda = 1 - \frac{p}{p_o}
$$
$$

Col 3: Plastic region ? (Y/N)  $\longrightarrow$  If  $\lambda > \lambda_{cr}$  then Y else N

Col 4:  $r_p / R \longrightarrow$  If  $\lambda < \lambda_{cr}$  (no plastic region) then  $r_p / R = 1$  else: If  $\varphi = 0:$   $\frac{r_p}{R} = \exp \left[ \frac{1}{2} (\lambda N_s - 1) \right]$ If  $\varphi > 0$ :  $\frac{r_p}{R} = \left\{ \left( \frac{2}{k+1} \right) \left[ \frac{2+N_s(k-1)}{2+N_s(k-1)(1-\lambda)} \right] \right\}^{\frac{1}{k-1}}$ 

**Si**  $\int$  $\sum_{i=1}^n a_i$  $\frac{P_o}{\sigma}$  $\left( 2G\right)$  $\left( p_{_{\alpha}}\right)$  $\left(\frac{10}{20}\right)$  $\sqrt{2G/2}$  $\binom{p}{2}$  $\left| \frac{P}{D} \right| \left| \frac{10}{2C} \right|$  $(R)(2G)$  $(r_n)(p_n)$  $=$  A  $_{cr}$   $\left(\frac{ }{R}\right)\left(\frac{ }{2G}\right)$ *p*<sup>0</sup> | *R* 1 2*G* 1 *r* 1 *n* 1 *R R* 1 2 (  $u_p$ ,  $v_p$   $p_o$ *cr p*  $2G$  ) and Col 5:  $u_p/R \longrightarrow \text{If } \lambda < \lambda_{cr}$  then  $u_p = n/a$  else:  $\frac{p}{R} = \lambda_{cr} \left| \frac{p}{R} \right| \left| \frac{P_o}{2C} \right|$ 

on of the convergence – confinement curve<br>
plastic region):  $\frac{u_R}{R} = \lambda \left(\frac{p_o}{2G}\right)$ <br>
else:  $\frac{u_R}{R} = \frac{u_p}{R} \left(\frac{r_p}{R}\right)^K$ <br>
for  $\lambda = 1$ <br>  $(u_{R\infty}/R)$ <br>  $\left[\frac{u_R}{u_{R\infty}}\right]^{0.588} - 1$ <br>
R) vs (p/p<sub>o</sub>), (x/R) vs (p/p<sub>o</sub>) o calculation of the convergence – confinement curve<br>  $\langle \lambda_{cr}$  (no plastic region):  $\frac{u_R}{R} = \lambda \left(\frac{p_o}{2G}\right)$ <br>
else:  $\frac{u_R}{R} = \frac{u_p}{R} \left(\frac{r_p}{R}\right)^K$ <br>
1 to  $u_R/R$  for  $\lambda = 1$ <br>  $(u_R/R) / (u_{R\alpha}/R)$ <br>  $= 1.10 \ln \left[\left(\frac{u_R}{u_{R\alpha}}\right)^$ calculation of the convergence – co<br>  $\lambda_{cr}$  (no plastic region):  $\frac{u_R}{R} = \lambda \left($ <br>
else:  $\frac{u_R}{R} = \frac{u_p}{R}$ <br>
to  $u_R/R$  for  $\lambda = 1$ <br>  $u_R/R$  /  $(u_{R\infty}/R)$ <br>
1.10  $\ln \left[ \left( \frac{u_R}{u_{R\infty}} \right)^{-0.588} - 1 \right]$ <br>
(b),  $(r_p/R) \text{ vs } (p/p_o)$ **he calculation of the conver**<br>  $\lambda < \lambda_{cr}$  (no plastic region)<br>
else:<br>
ual to  $u_R/R$  for  $\lambda=1$ <br>  $\frac{u_R}{R} = 1.10 \ln \left[ \left( \frac{u_R}{u_{R\infty}} \right)^{-0.588} -1 \right]$ <br>  $(p/p_0)$ ,  $(r_p/R)$  vs  $(p/p_0)$ , (GRC) EXCEL spreadsheet for the calculation of the convergence – confinement curve  $= \lambda \left(\frac{p_o}{2G}\right)$  $u_R$  **o**  $\begin{bmatrix} p_o \end{bmatrix}$ Col 6:  $u_R / R \longrightarrow$  If  $\lambda < \lambda_{cr}$  (no plastic region):  $\frac{u_R}{R} = \lambda \left( \frac{p_o}{2G} \right)$  $\lambda$  |  $\frac{P_o}{\sim}$  |  $\frac{1}{\sim}$  $R$  *G*  $\setminus$  *2G*  $\setminus$  *2G*  $\setminus$  *3G*  $\setminus$  *3G*  $\setminus$  *3G*  $\setminus$  *3G*  $\setminus$  *5G*  $\setminus$ *K*  $\left(r_{n}\right)^{\Lambda}$  $\bigwedge$  and  $\bigwedge$  and  $\bigwedge$ *u r*  $u_n$ ,  $u_n$ ,  $u_n$ ,  $u_n$ else:  $\frac{R}{R} = \frac{p}{R} \left| \frac{p}{R} \right|$  $\begin{array}{|c|c|c|c|c|}\n\hline\nD & \multicolumn{1}{|c|}{c} \\
\hline\nD & \multicolumn{1}{|c|}{c} \\
\hline\n\end{array}$ **Contract Contract**  $\equiv$  $(R)$  $\int$  and  $\int$  and  $\int$ *R R R* Calculate  $u_{R\infty}/R$ : equal to  $u_R/R$  for  $\lambda=1$ Col 7:  $u_R / u_{R\infty} \longrightarrow (u_R / R) / (u_{R\infty} / R)$  $-0.588$  $0.588$  $x \left[ \begin{array}{c} u_R \end{array} \right]$   $\left[ \begin{array}{c} u_R \end{array} \right]$ Col 8: x / R  $\infty$  , we have  $\sim$  $R$ ∞  $\;$   $\;$ Plot curves:  $(u_R / R)$  vs  $(p/p_o)$ ,  $(r_p / R)$  vs  $(p/p_o)$ ,  $(x/R)$  vs  $(p/p_o)$  or  $(u_R / R)$ Ground Reaction (GRC) Longitudinal Diplacement (LDP)  $u_{Roe} = \frac{p_o}{2G} R$ Limit of elastic behaviour  $c_u = \frac{c \cos\varphi}{1 - \sin\varphi}$ Mote: Curve III does

 $\overline{u_R}$ 

Elastic ground

 $\lambda$  = Deconfinement = 1 -  $p/p_a$ 

EXCEL spreadsheet for the calculation of the convergence – confinement curve

else:

Col 9: 
$$
\sigma_r / p_o \text{ (at r=R): } \longrightarrow \frac{\sigma_r}{p_o} = (1 - \lambda)
$$

Col 10:  $\sigma_{\theta}$ / $p_o$  (at r=R):  $\longrightarrow$  If  $\lambda < \lambda_{cr}$  then:  $\frac{\nu_{\theta}}{\nu} = (1 + \lambda)$  $p_{o}$  $\frac{\sigma_{\theta}}{2} = (1 + \lambda)$ 

Plot curves: ( $\sigma_r / p_o \& \sigma_{\theta} / p_o$ ) vs ( $p/p_o$ )





 $\begin{array}{c} \mathbf{r} \\ \mathbf{r} \end{array}$  2 *r* 

 $p_{o}$   $\left(p_{o}^{+}\right)$   $N_{s}$ 

 $\left(\frac{r}{p_o}\right) + \frac{r}{N_s}$ 

*o o s*

 $k \rvert$   $\rvert$   $\rvert$   $\rvert$   $\rvert$   $\rvert$   $\rvert$   $\rvert$   $\rvert$ 

 $\sigma_{\theta} = k \left( \sigma_r \right) + \frac{2}{k}$ 

### Examples of Panet – Chern curves :



I  $\setminus$  $\bigg($  $\lambda = f\left| \frac{\mu}{\sigma} \right|; N_s, \varphi, \delta$ *R x f*

▎

 $\int$ 

#### Examples of Panet – Chern curves :





### Examples of radius of plastic zone:

$$
\frac{r_{P}}{R}=f\left(\lambda \ ; \ N_{s}, \phi\right)
$$





### Examples of tunnel wall displacement:





Example: Convergence – confinement curve  $(u_R)$  -  $(\lambda)$ 

$$
\frac{u_R}{R} = f\left(\lambda, \frac{p_o}{2G} \ ; \ N_s, \varphi, \delta\right) \qquad \frac{u_R}{u_R}
$$

$$
\frac{u_R}{u_{R\infty}} = f(\lambda \ ; \ N_s, \varphi, \delta)
$$



### Example: Radius of the plastic zone



г,

R



# Example: Radius of plastic zone





### Example: Panet - Chern curve



Example: wall displacement curve (A)

$$
\frac{u_R}{u_{R\infty}} = f\left(\frac{x}{R}\right) \qquad \text{(Chern)}
$$

$$
\frac{u_R(x)}{u_{R\infty}} = \left[1 + exp\left(0.91\frac{x}{R}\right)\right]^{-1.7}
$$



# Example: wall displacement curve (B)



